ADOMIAN DECOMPOSITION METHOD WITH NEUMANN BOUNDARY CONDITIONS FOR SOLUTION OF NONLINEAR BOUNDARY VALUE PROBLEM

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Abstract: The Adomian decomposition method (ADM) is a creative and effective method for exact solution of functional equations of various kinds. Adomian decomposition method solves wide class of linear and non-linear, ordinary or partial differential equations. This paper presents the Adomian decomposition method for the solution of nonlinear boundary value problem using Neumann boundary conditions. In this approach, the solution is found in the form of a convergent power series with easily computed components. To show the efficiency of the method, numerical results and graphical representation of results are presented and compared with exact solution.

Keywords: Adomian decomposition method, Neumann boundary conditions, Nonlinear boundary value problem, Bratu problem, Burger problem.

1. INTRODUCTION

The Adomian Decomposition Method [1-3] has been utilized for the solution of linear and nonlinear ordinary and partial differential equations. Some authors known as Kaya [4], El-Saved, Biazar Hashim, and Lesnic [5-7] examined some different scientific models analytically and numerically. Another two writers named Sweilam and Khader [8-10] used the ADM for the analysis of nonlinear atmosphere of multiwalled carbon Nano tubes. ADM moreover gives a convergent sequence of estimations for some terms of high accuracy. Cherruault [11, 12] discussed the convergence of the ADM. This method was proved very valuable for solving boundary value problems by many other writers. There was a reason of its value, which was shortest use of restraining supposition, linearization or Green functions. Rach and Adomian [13, 14] exemplified the way to solve nonlinear BVPs in only some measurement with the help of decomposition technique and considered many ordinary and partial differential equations through the two conditions known as Dirichlet and Neumann boundary [15].

2. MATERIAL AND METHODS

Let the two-point boundary value problem be u''(x) + p(x)u'(x) + q(x)f(u(x)) = r(x) (1) $x \in [a, b],$ Along with the Neumann boundary conditions $u'(a) = \alpha, u'(b) = \beta$ (2) Suppose the inverse operator $L_{xx}^{-1}[.] = \iint [.] dx dx$

is relate to Eq. (1) resultant in $u(x) = \varphi_x + L_{xx}^{-1}(r(x) - p(x)u'(x) - q(x)f(u(x))(3)$ where $\varphi_x = c_1 + c_2x$, c_1 and c_2 are constants, the linear and nonlinear terms u and f(u(x)) as follows:

$$\varphi_{x} = \sum_{n=0}^{\infty} \varphi_{x,n} = \sum_{n=0}^{\infty} (c_{1,n} + c_{2,n}x),
u(x) = \sum_{n=0}^{\infty} u_{n}(x),
f(u(x)) = \sum_{n=0}^{\infty} A_{n}(u_{0}, u_{1}, \cdots, u_{n})\}$$
(4)

where A_n the Adomian polynomials that are computed as follows

$$A_n = \frac{1}{n!} \left[\frac{d^n}{d\lambda^n} f(\sum_{i=0}^{\infty} \lambda^i u_i) \right]_{\lambda=0}, \ n \ge 0$$

The decomposition of the initial expression is required substitute (4) into (3),

$$\begin{split} u_0(x) &= \varphi_{x,0} + L_{xx}^{-1} r(x) \\ u_{n+1}(x) &= \varphi_{x,n} - L_{xx}^{-1} (p(x) u'_n(x) + q(x) A_n), \quad n \geq 0 \\ \text{The Eq. (1) is given by} \\ \varphi_m(x) &= \sum_{n=0}^{\infty} u_n(x) \text{ for } m > 0 \\ \text{Thus } \varphi_1 &= u_0, \varphi_2 = \varphi_1 + u_1, \varphi_3 = \varphi_2 + u_2 \text{ , etc.} \\ \varphi_1 &= u_0 = c_{1,0} + c_{2,0} x + L_{xx}^{-1} r(x). \\ \text{Using (2) results in} \\ c_{2,0} + [L_{xx}^{-1} r(x)]'_{x=a} = \alpha \\ c_{2,0} + [L_{xx}^{-1} r(x)]'_{x=b} = \beta \\ \text{If } u'(a) &= \alpha \text{ and } u'(b) = \beta \text{ are ordinary differential equation of } \end{split}$$

second order through Neumann boundary conditions

$$L^{-1}u''(x) = u(x) - (x - \Omega)u'(\alpha) - \frac{\Omega}{2}u'(b) - \frac{\Omega}{2}$$

$$\frac{1}{a} \int_{0}^{a} u(x) dx , \quad a \le x \le b$$
where $L^{-1}_{-1}[.]$ is obvious by means of

$$L_{xx}^{-1}[.] = \int_{\Omega}^{x} dx' \int_{a}^{x'} [.] dx'' + \frac{1}{\alpha} \int_{0}^{\Omega} dx' (x' \int_{b}^{x'} [.] dx'') (5)$$

Where Ω is random
$$L_{xx}^{-1}[.] = \int_{\Omega}^{x} dx' \int_{a}^{x'} [.] dx'' + \frac{1}{\alpha} \int_{0}^{\Omega} dx' (x' \int_{b}^{x'} [.] dx'')$$

Then
$$L_{xx}^{-1}[u''(x)] = \int_{\Omega}^{x} dx' \int_{a}^{x'} [u''(x)] dx'' + \frac{1}{\alpha} \int_{0}^{\Omega} dx' (x' \int_{a}^{x'} [u''(x)] dx'' + \frac{1}{\alpha} \int_{0}^{\Omega} (x' \int_{a}^{X'} [u''(x)] dx'' + \frac{1}{\alpha} \int_{0}^{\Omega} (x' \int_{a}^{X'} [u''(x)] dx'' + \frac{1}{\alpha} \int_{0}^{X'} [u''(x)] dx'' + \frac{1}{\alpha} \int_{0}^{\Omega} (x' \int_{a}^{X'} [u''(x)] dx'' + \frac{1}{\alpha} \int_{0}^{X'} [u''(x)] dx'' + \frac{1}{\alpha} \int_{0}^{$$

$$\begin{aligned} & \frac{1}{\alpha} \int_{0}^{\alpha} dx (x \int_{b}^{\alpha} [u(x)]dx) \\ &= \int_{\alpha}^{x} [u'(x') - u'(a)]dx' + \frac{1}{\alpha} \int_{0}^{\alpha} x'[u'(x') - u'(b)]dx' \\ &= u(x) - u(\Omega) - (x - \Omega)u'(a) + \frac{1}{\alpha} \int_{0}^{\alpha} x'u'(x')dx' - \\ & \frac{1}{\alpha} \left[\frac{\alpha^{2}}{2} u'(b) \right] \\ &= u(x) - u(\Omega) - (x - \Omega)u'(a) + \frac{1}{\alpha} \int_{0}^{\alpha} x'u'(x')dx' - \\ & \frac{\alpha}{2} u'(b) \\ &= u(x) - u(\Omega) - (x - \Omega)u'(a) + \frac{1}{\alpha} [\Omega u(\Omega) - \\ & \int_{0}^{\alpha} u(x')dx'] - \frac{\alpha}{2} u'(b) \\ &= \\ & u(x) - u(\Omega) - (x - \Omega)u'(a) + [u(\Omega) - \\ & \frac{1}{\alpha} \int_{0}^{\alpha} u(x')dx'] - \frac{\alpha}{2} u'(b) \end{aligned}$$

$$\frac{1}{\alpha} \int_0^{\alpha} u(x') dx' - \frac{1}{2} u'(b) = u(x) - (x - \Omega)u'(a) - \frac{\alpha}{2} u'(b) - \frac{1}{\alpha} \int_0^{\Omega} u(x) dx$$

There is a setup of algorithms used for linear and nonlinear second-order ordinary and partial differential equations by the help of Neumann boundary form.

u''(x) = r(x) - p(x)u'(x) - q(x)f(u(x))Now, the operator $L_{xx}^{-1}(.)$ by Eq. (5), it will help to obtain $u(x) = (x - \Omega)u'(a) + \frac{\Omega}{2}u'(b) + \frac{1}{\alpha}\int_{0}^{\Omega}u(x)dx + L_{xx}^{-1}(u(x))u'(x) = \sigma(u)u'(x)$

$$L_{xx}^{-1}(r(x) - p(x)u'(x) - q(x)f(u(x)))$$

When u and f (u(x)) are decomposed by Eq. (4)
$$u_{0} = (x - \Omega)u'(a) + \frac{\alpha}{2}u'(b) + L_{xx}^{-1}[r(x)]$$
$$u_{n+1} = \frac{1}{\alpha}\int_{0}^{\alpha}u_{n}(x)dx - L_{xx}^{-1}[p(x)u'_{n}(x) + q(x)A_{0}$$
as well as within the linear form
$$u_{0} = (x - \Omega)u'(a) + \frac{\alpha}{2}u'(b) + u_{xx}^{-1}[r(x)]$$
$$u_{n+1} = \frac{1}{\alpha}\int_{0}^{\alpha}u_{n}(x)dx - L_{xx}^{-1}[p(x)u'_{n}(x) + q(x)u_{n}] , n \ge 0$$

The estimated analytic explanation is given by $\varphi_n(x;\Omega) = \sum_{i=0}^{\infty} u_i(x;\Omega)$

3. NUMERICAL ILLUSTRATIONS

3.1 Bratu Problem

The following Bratu problem $u^{\prime\prime}(x)-2e^u=0\,,$ $0 \le x \le 1$ u'(0) = 0, $u'(1) = 2 \tan(1)$ Application of ADM gives $u_0 = (x - \Omega)u'(a) + \frac{\alpha}{2}u'(b) + L_{xx}^{-1}[r(x)]$ $u_0 = \Omega \tan(1)$ $u_{n+1} = \frac{1}{a} \int_0^{a} u_n(x) dx + \int_a^{x} dx' \int_0^{x'} [2A_n(x'')] dx''$ $+\frac{1}{o}\int_{0}^{a}dx'(x'\int_{1}^{x'}[2A_{n}(x'')]dx'')$

where $n \ge 0$ and A_n be the Adomian polynomials which are $A_0 = e^{u_0}$

$$A_{1} = u_{1}e^{u_{0}}$$
$$A_{2} = \left(u_{2} + \frac{u_{1}^{2}}{2}\right)e^{u_{0}}$$

 $A_3 = \left(u_3 + u_1 u_2 + \frac{u_1^3}{6}\right) e^{u_0}$ The estimated solutions $\Phi 2(x)$, $\Phi 3(x)$ and $\Phi 4(x)$ as $\Omega \rightarrow 0$ are given by

$$u_{1} = \frac{1}{\Omega} \Omega \tan(1)(\Omega) + \int_{\Omega}^{x} dx' [2e^{u_{0}}] dx'' + 0$$

= $\Omega \tan(1) + \int_{\Omega}^{x} 2e^{\Omega \tan(1)} dx''$
= $\Omega \tan(1) + 2[e^{\Omega \tan(1)}x]_{\Omega}^{x}$
= $\Omega \tan(1) + 2e^{\Omega \tan(1)}\left(\frac{x^{2}}{2}\right)$
 $u_{1} = x^{2}$
 $\varphi_{2}(x) = u_{0} + u_{1}, \ \varphi_{2}(x) = \Omega \tan(1) + x^{2}$
 $\varphi_{2}(x) = x^{2}$
 $\varphi_{3}(x) = x^{2} + \frac{x^{4}}{6}$
 $\varphi_{4}(x) = x^{2} + \frac{x^{4}}{6} + \frac{2x^{6}}{45}$

The exact solution is $u(x) = -2 \log [\cos(x)]$. The Maclaurin series of this solution is

			TI		•	E		Б	•	
	Tab	le-1: \$	Sum	mar	v of re	sults fr	om	Bratu	proble	m
и(л) – .	λ Τ (5	45	1260	14175	Τ	•		
11(v) — ·	$r^2 \pm \frac{x}{2}$	4 	$2x^{6}$	$\perp \frac{17x^8}{17x^8}$	$+ \frac{62x^{10}}{10}$	⊥			

х	U	Exact	Error
0	0	0	0

0.05	.002501041667194	.002501042361638	0.00000000694444
0.1	.010016666802025	.01006711246470	0.000000044444446
0.15	.02258437843090	.022584884733284	0.000000506250193
0.2	.040266701654236	.040269546104816	0.000002844450580
0.25	.063151251710491	.063162102494940	0.000010850784448
0.3	.091350911041714	.091383311852116	0.000032400810402
0.35	.125004200567281	.125085906475610	0.000081705908329
0.4	.164275967461813	.164458038150110	0.000182070688298
0.45	.209358551413617	.209727717240980	0.000369165827363
0.5	.260473641424162	.261168480887446	0.000694839463283
0.55	.317875094898601	.319106610717962	0.000694839463283
0.6	.38185306214400	.383930338838876	0.002077276694875
0.65	.452739843933609	.45610648239232	0.003361804305624
0.7	.530918009062173	.536171515135862	0.005253506073689
0.75	.616831414358957	.624799795798210	0.007968381439254
0.8	.710999902991803	.722781493622688	0.011781590630885
0.85	.814038609637238	.83108196165270	0.017043352014031
0.9	.926682974758286	.950884887171628	0.024201912413343
1	1.18453262786596	1.23125940772028	0.046720312906067





Suppose the nonlinear Burger problem

$$u'' + uu' + u = \frac{1}{2}\sin(2x), \quad 0 \le x \le \frac{\pi}{2}$$
$$u'(0) = 1, \ u'\left(\frac{\pi}{2}\right) = 0$$

Consider the ADM; the solution can be computed as follows $u_0 = (x - \Omega)u'(a) + \frac{\alpha}{2}u'(b) + L_{xx}^{-1}[r(x)]$ = $(x - \Omega) + L_{xx}^{-1}[\frac{1}{2}\sin(2x)]$ $= (\mathbf{x} - \Omega) + \int_{\Omega}^{x} dx' \int_{0}^{x'} [\frac{1}{2} \sin(2x)] dx'' +$ $\frac{1}{\Omega} \int_{0}^{\Omega} dx' (x' \int_{\frac{\pi}{2}}^{x'} \frac{1}{2} \sin(2x)) dx''$

$$= (x - \Omega + \int_{\Omega}^{x} dx' \left| \frac{1}{2} \cdot \frac{-\cos(2x)}{2} \right|_{0}^{x'} + \frac{1}{\Omega} \int_{0}^{\Omega} dx' \left(x' \left| \frac{1}{2} \cdot \frac{-\cos(2x)}{2} \right| \frac{x'}{2} \right)$$

	0	1011000	
0	0	0	0
.05	.049979169270678	.4997916927068	0.00000000069389
.1	.099833416646828	.099833416646828	0.00000000138778
.15	.14938132473599	.149438132473599	0.00000000555112
.2	.198669330795062	.198669330795061	0.00000005273559
.25	.247403959254529	.247403959254523	0.000000059674488
.3	.295520206661384	.295520206661340	0.000000059674488
.35	.342897807455693	.342897807455451	0.000002416400413





The following nonlinear BVP $u'' - (u')^2 = 0$ u'(0) = -1, $u'(1) = \frac{-1}{2}$ We obtain the following $A_n = (n+1)\sum_{i=1}^n u'_{n-i}u'_i$ $u_0 = (x - \Omega)u'(a) + \frac{\Omega}{2}u'(b) + L_{xx}^{-1}[r(x)]$ $= -x + \Omega - \frac{\Omega}{4}$ $u_0 = -x + \frac{3\Omega}{4}$, $u_0 = -x + \mu_0$ $u_0 = -x + \mu_0$ $u_{n+1} = \int_{\Omega}^x (\int_0^x A_n dx) dx + \frac{1}{\Omega} \int_0^\Omega (x \int_1^x A_n dx) dx + \frac{1}{\Omega} \int_0^\Omega u_n dx$

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The following approximants

$$u_{1} = \int_{\Omega}^{x} \int_{0}^{x} A_{0} dx + \frac{1}{\alpha} \int_{0}^{\Omega} (x \int_{1}^{x} A_{0} dx + \frac{1}{\alpha} \int_{0}^{\Omega} u_{0} dx)$$

$$= \int_{\Omega}^{x} \int_{0}^{x} (u_{0}')^{2} dx \quad , \int_{\Omega}^{x} \int_{0}^{x} dx dx$$

$$u_{1} = \frac{x^{2}}{2} - \frac{\Omega^{2}}{2}$$

$$u_{1} = \frac{x^{2}}{2} + \mu_{1}$$

$$\varphi_{2}(x) = u_{0} + u_{1}$$

$$\varphi_{2}(x) = -x + \frac{x^{2}}{2} + \mu_{1}$$

$$\varphi_{2}(x) = -x + \frac{x^{2}}{2} + \mu_{1}$$

$$\varphi_{3}(x) = -x + \frac{x^{2}}{2} - \frac{x^{3}}{3} + \mu_{2}$$

$$\varphi_{4}(x) = -x + \frac{x^{2}}{2} - \frac{x^{3}}{3} + \frac{x^{4}}{4} - \frac{x^{5}}{5} + \mu_{4}$$

where μ_{1} are constants in terms of Ω . As $n \to \infty$.

where	μ_i are	constants	in ten	115 01 22	. As n	$, \infty, \text{ men } u$	_
-log(x-	+1) +µ,	where µ=	limit µ	ι _i as n —	→ ∞		

-	Table 5. Summary of results from above nominear DVI							
X	u	Exact	Error					
0	0	0	0					
0.05	-0.0487901666666667	-0.048790164169432	0.00000002497235					
0.1	-0.095310333333333	-0.095310179804325	0.000000153529008					
0.15	-0.139763625000000	-0.139761942375159	0.000001682624841					
0.2	-0.1823306666666667	-0.182321556793955	0.000009109872712					
0.25	-0.223177083333333	-0.223143551314210	0.000033532019124					
0.3	-0.26246100000000	-0.262364264467491	0.000096735532509					
0.35	-0.300340541666667	-0.300104592450338	0.000235949216329					
0.4	-0.336981333333333	-0.336472236621213	0.000509096712120					
0.45	-0.37256400000000	-0.371563556432483	0.001000443567517					
0.5	-0.4072916666666667	-0.405465108108164	0.001826558558502					
0.55	-0.441397458333333	-0.438254930931155	0.003142527402178					
0.6	-0.47515200000000	-0.470003629245736	0.005148370754264					

Table 3: Summary of results from above nonlinear BVP







Figure 6: Graph for above nonlinear BVP with iterations 3.4 Nonhomogeneous Wave Equation Let the nonhomogeneous wave equation be $u_{tt} = u_{xx} + 2\pi^2 e^{-\pi t} \sin(\pi x),$ $0 \le x \le 1$, $t \ge 0$ $u_x(0,t) = \pi e^{-\pi t}, \ u_x(1,t) = -\pi e^{-\pi t}$ We use the recursion method $u_0 = (x - \Omega)u'(a) + \frac{\Omega}{2}u'(b) + L_{xx}^{-1}[r(x)]$ = $(x - \Omega)\pi e^{-\pi t} + \frac{\Omega}{2}(-\pi e^{-\pi t}) - L_{xx}^{-1}[2\pi^2 e^{-\pi t}\sin(\pi x)]$ $(x-\Omega)\pi e^{-\pi t} + \frac{\Omega}{2}(-\pi e^{-\pi t}) +$ $2\pi^{2}e^{-\pi t} \left[\int_{a}^{x} dx' \int_{0}^{x'} \sin(\pi x) dx'' + \frac{1}{a} \int_{0}^{a} dx' (x' \int_{1}^{x'} \sin(\pi x)) dx'''\right]$ $(x-\Omega)\pi e^{-\pi t} + \frac{\Omega}{2}(-\pi e^{-\pi t}) +$ $\int_{\Omega}^{x} dx' \left| \frac{-\cos(\pi x)}{\pi} \right|_{0}^{x'} + \frac{1}{\Omega} \int_{0}^{\Omega} dx' \left(x' \left| \frac{-\cos(\pi x)}{\pi} \right|_{1}^{x'} \right)$ $= (x - \Omega)\pi e^{-\pi t} + \frac{\Omega}{2}(-\pi e^{-\pi t}) + \int_{\Omega}^{x} dx' \left(\frac{-\cos(\pi x')}{\pi} + \frac{1}{\pi}\right) +$ $\frac{1}{\Omega}\int_0^\Omega dx' \left(x' \left(\frac{-\cos(\pi x')}{\pi} - \frac{1}{\pi}\right)\right)$ $= (x - \Omega)\pi e^{-\pi t} + \frac{\alpha}{2}(-\pi e^{-\pi t}) + \left|\frac{-\sin(\pi x')}{\pi^2}\right|_{\Omega}^{X} + \frac{1}{\pi}(x - \Omega) + \frac{1}{\alpha}\left|x' \cdot \frac{-\sin(\pi x')}{\pi^2}\right|_{\Omega}^{\Omega} - \int_{0}^{\Omega} \frac{-\sin(\pi x')}{\pi^2} dx' - \frac{1}{\pi}\left(\frac{\alpha^2}{2}\right)$ $= (x - \Omega)\pi e^{-\pi t} + \frac{\alpha}{2}(-\pi e^{-\pi t}) + \frac{-\sin(\pi x)}{\pi^2} + \frac{\sin(\pi \Omega)}{\pi^2} + \frac{x}{\pi} - \frac{\alpha}{\pi} + \frac{1}{\alpha}\left[-\Omega\frac{\sin(\pi \Omega)}{\pi^2} + \left|\frac{-\cos(\pi x')}{\pi^3}\right|_{\Omega}^{\Omega} - \frac{\alpha^2}{2\pi}\right]$ $= (x - \Omega)\pi e^{-\pi t} + \frac{\alpha}{2}(-\pi e^{-\pi t}) + \frac{-\sin(\pi x)}{\pi^2} + \frac{\sin(\pi \Omega)}{\pi^2} + \frac{x}{\pi} - \frac{\alpha}{\pi} - \frac{\sin(\pi \Omega)}{\pi^2} - \frac{\cos(\pi \Omega)}{\pi^2} + \frac{1}{\pi} - \frac{\alpha^2}{\pi^2}$ $= (x - \Omega)\pi e^{-\pi t} + \frac{1}{2}(-\pi e^{-\pi t}) + \frac{\pi^2}{2\Omega\pi} + \frac{\pi^2}{\pi^2} + \frac{\pi^2}{\pi^2}$ $\frac{\cos(\pi\Omega)}{\Omega\pi^3} + \frac{1}{\Omega\pi^3}$ = $(x-\Omega)\pi e^{-\pi t}+\frac{\Omega}{2}(-\pi e^{-\pi t})+$ $2\pi^{2}e^{-\pi t}\left(\frac{-\sin(\pi x)}{\pi^{2}}\right) + \frac{2\pi^{2}e^{-\pi t}}{\pi}x - 2\pi^{2}e^{-\pi t}\frac{3\Omega}{2\pi} - 2\pi^{2}e^{-\pi t}\frac{\cos(\pi\Omega)}{\Omega\pi^{3}} + 2\pi^{2}e^{-\pi t}\cdot\frac{1}{\Omega\pi^{3}}$



$$\begin{split} u_{n+1} &= \frac{1}{\alpha} \int_0^{\alpha} u_n(x,t) dx + \int_{\alpha}^{x} dx' \int_0^{x'} \frac{\partial^2 u_n(x'',t)}{\partial t^2} dx'' + \\ & \frac{1}{\alpha} \int_0^{\alpha} dx'(x' \int_1^{x'} \frac{\partial^2 u_n(x'',t)}{\partial t^2} dx'') \\ \varphi_5(x,t) &= e^{-\pi t} \left[\pi x - \frac{(\pi x)^3}{3!} + \frac{(\pi x)^5}{5!} - \frac{(\pi x)^7}{7!} + \frac{(\pi x)^9}{9!} - \\ & \frac{(\pi x)^{11}}{19958400} + \cdots \right] \end{split}$$

Estimate more terms therefore

$$u(x,t) = \lim_{n \to \infty} \varphi_n(x,t) = e^{-\pi t} \sin(\pi x)$$

Which matches the accurate solution

$$u(x,t) = e^{-\pi t} \sin(\pi x) + \mu_5 t + \mu_5 t$$

where μ_5 and μ_6 are arbitrary therefore μ_5 and μ_6 vanish. Table 4: Summary of results obtained in example 4

х	t	u	Exact	Error(1.0e-03)
0	0	0	0	0
0.05	0.05	0.133694525331659	0.133694525331659	0.000000000000056
0.1	0.1	0.225706844271172	0.225706844271226	0.00000000053818
0.15	0.15	0.283393778539763	0.283393778543749	0.00000003985257
0.2	0.2	0.313576432136283	0.313576432217014	0.000000080731533
0.25	0.25	0.322396941140479	0.322396941944834	0.000000804354972
0.3	0.3	0.315242477067670	0.315242482184127	0.000005116456370
0.35	0.35	0.296721574800026	0.270679760789519	0.000023880498135
0.4	0.4	0.270679671921842	0.296721598680524	0.000088867676273
0.45	0.45	0.240242625238963	0.240242903426645	0.000278187682112
0.5	0.5	0.207878816529668	0.207879576350762	0.000759821093571
0.55	0.55	0.175472204534316	0.175474063168161	0.001858633845164
0.6	0.6	0.144400277908527	0.144404428880597	0.004150972069766
0.65	0.65	0.115612286868291	0.115620875625684	0.008588757393277
0.7	0.7	0.089704366707892	0.089721018894541	0.016652186648619
0.75	0.75	0.066989212680271	0.067019739708273	0.030527028002325
0.8	0.8	0.047558829607868	0.047612129067901	0.053299460032789
0.85	0.85	0.031339573569087	0.031428732320961	0.089158751874673
0.9	0.9	0.018139243360200	0.018282839453758	0.143596093558532
0.95	0.95	0.007686384523204	0.007909971253941	0.223586730737733
1	1	0.000337742318605	0.0000000000000000000000000000000000000	0.337742318605137







Figure 8: Graph of exact solution for nonhomogeneous wave equation in 3-D

4. CONCLUSION

Adomian decomposition method has been known to be a powerful device for solving many functional equations as algebraic equations, ordinary and partial differential equations, integral equations and so on. Here we used this method for solving nonlinear BVP. It is demonstrated that this method has the ability of solving systems of both linear and non-linear differential equations. In above problems, there was a nonlinear system and we derived the exact solutions. For non-linear systems, we usually derive a very good approximation to the solutions with the Neumann boundary conditions. It is also important that the Adomian decomposition method does not require discretization of the variables. It is not affected by computation round errors and one is not faced with necessity of large computer memory and time. Comparing the results with exact solutions, the Adomian decomposition method was clearly reliable techniques. It is important that this method unlike the most numerical techniques provides a closed form of the solution.

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